

X-ray visualized interfaces in high-speed sprays

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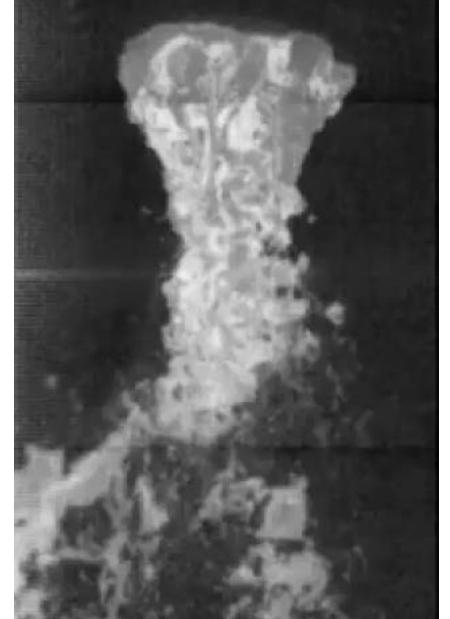
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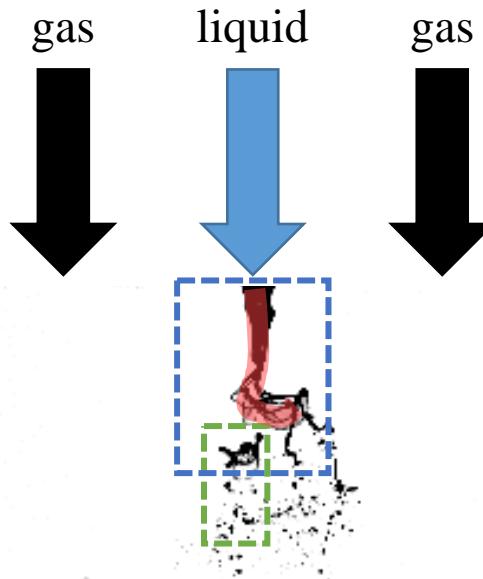


- Liquid-gas flows are critical in engineering process innovation and intensification
- Liquid sprays are critical for combustion systems, manufacturing, heat management, chemical processing, painting, e. g.:
 - Liquid fuel sprays
 - Liquid metal atomization
 - Spray cooling and coating
 - Pharmaceutical, food, consumer products
 - Fire safety
 - Ship wake and sea spray



Fe-Cr ODS Alloy at ~1800 C
I. Anderson, Ames Lab.





Assisted atomization: breaking of a liquid jet into a spray (droplet cloud) by a gas co-flow

Spray formation:

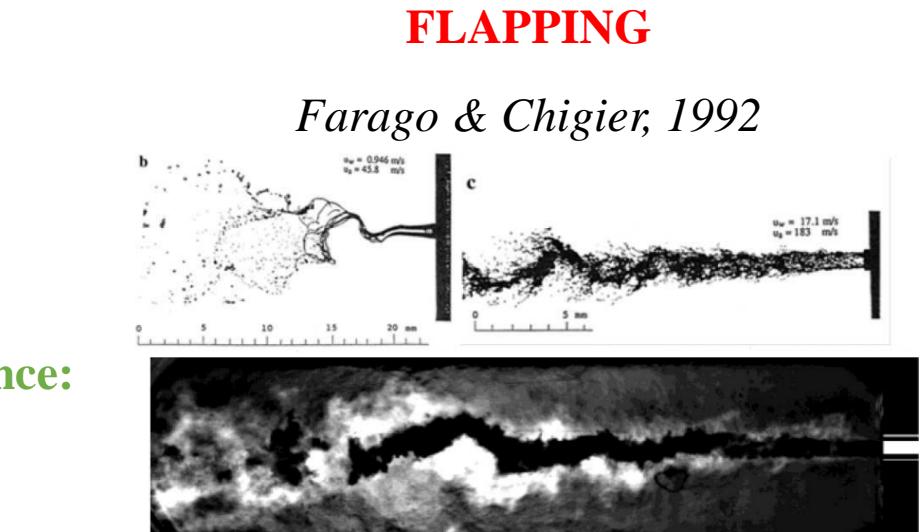
- Interfacial instabilities
- Primary break-up

Drops/ligaments in turbulence:

- Secondary break-up
- Turbulent dispersion

Droplets in turbulence:

- Turbulent dispersion
- Evaporation



Cryogenic, *Locke et al. 2010*

- Frequency well modeled (Delon et al., 2018)
- Flapping affects the cascade of mechanisms, up to droplet spatiotemporal distributions
- Dimensionality and role of swirl



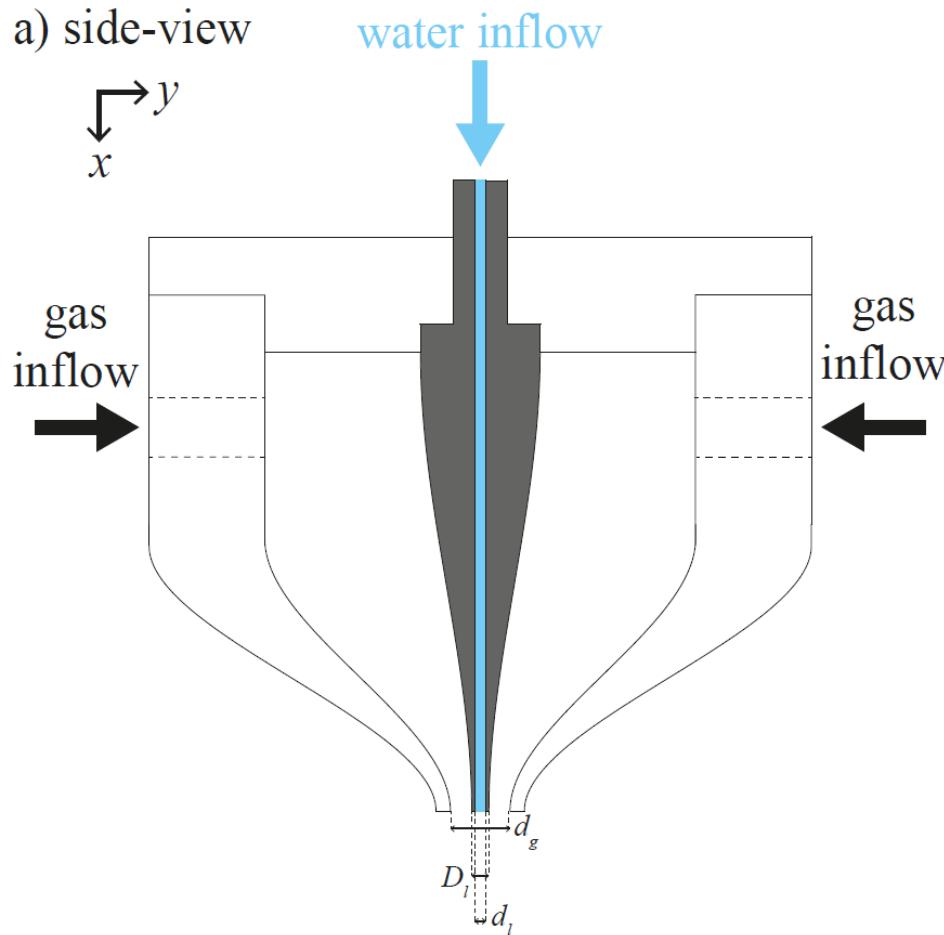
Effect for high-speed sprays

O. Desjardins' group

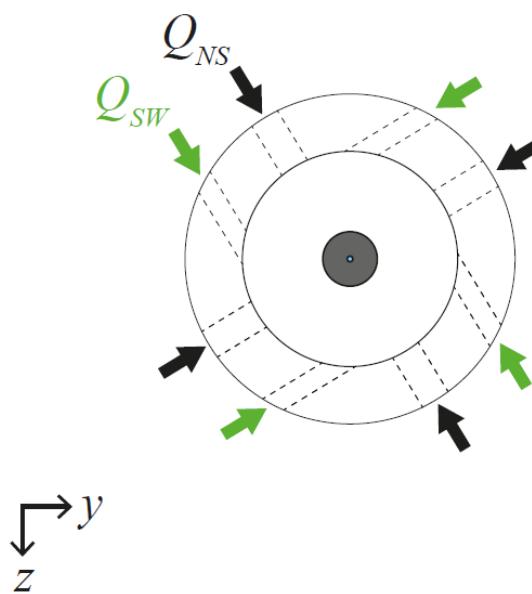
L. Vu et al., IJMF 2023

$$Q_{Total} = Q_{SW} + Q_{NS} = cst$$

a) side-view



b) top-view



$$SR = \frac{Q_{SW}}{Q_{NS}}$$



Threshold of angular to longitudinal momenta ratio



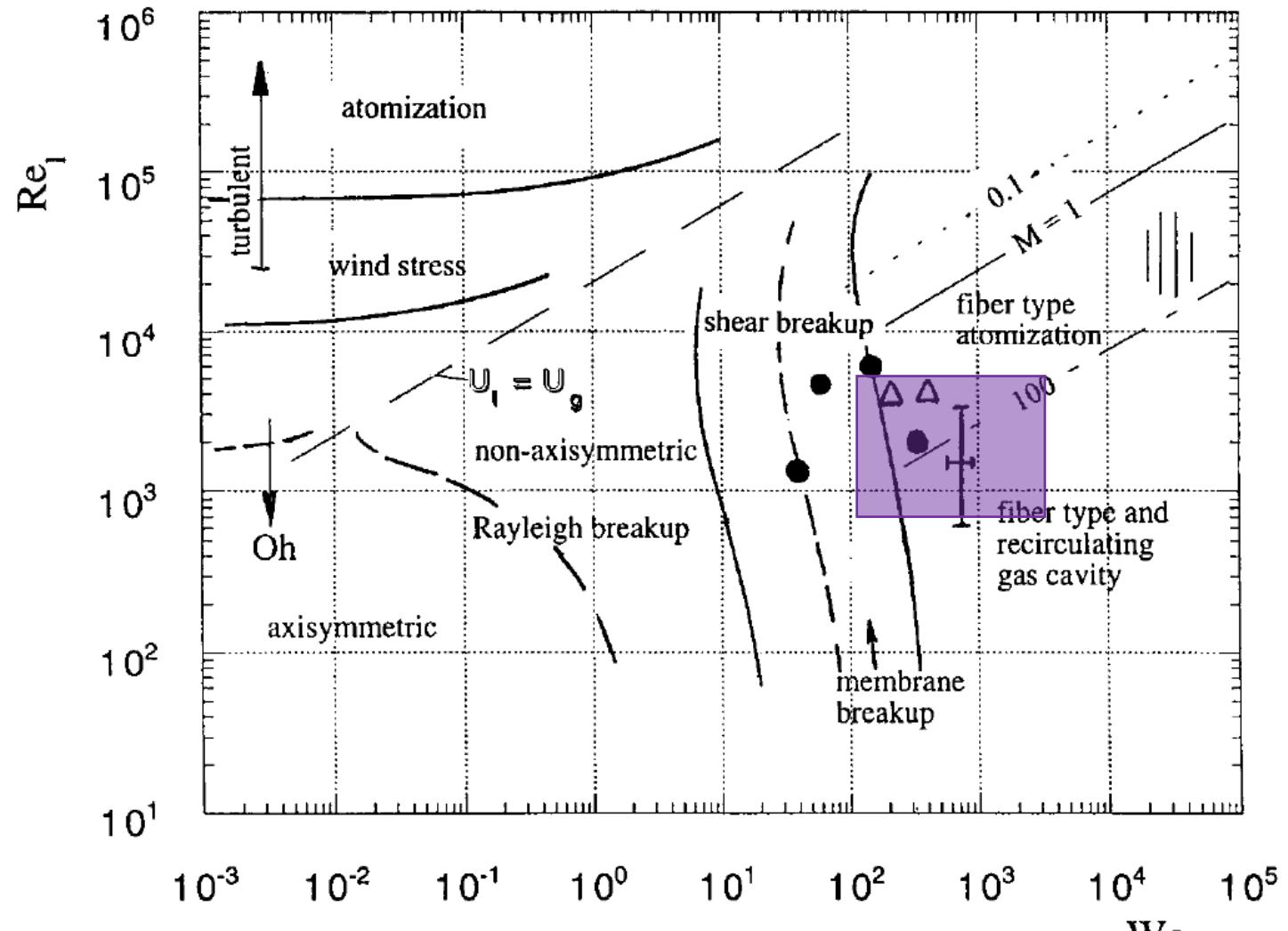
Lasheras & Hopfinger
ARFM 2000

Control parameters: liquid flow rate, no-swirl and swirl gas flow rates $\rightarrow Re_l, Re_g, SR$

$$\text{Weber number } We_g = \frac{\rho_g u_g^2 d_l}{\sigma}$$

$$\text{Gas-to-liquid dynamic pressure ratio } M = \frac{\rho_g u_g^2}{\rho_l u_l^2}$$

Qualitative spray formation pictures in the near-field

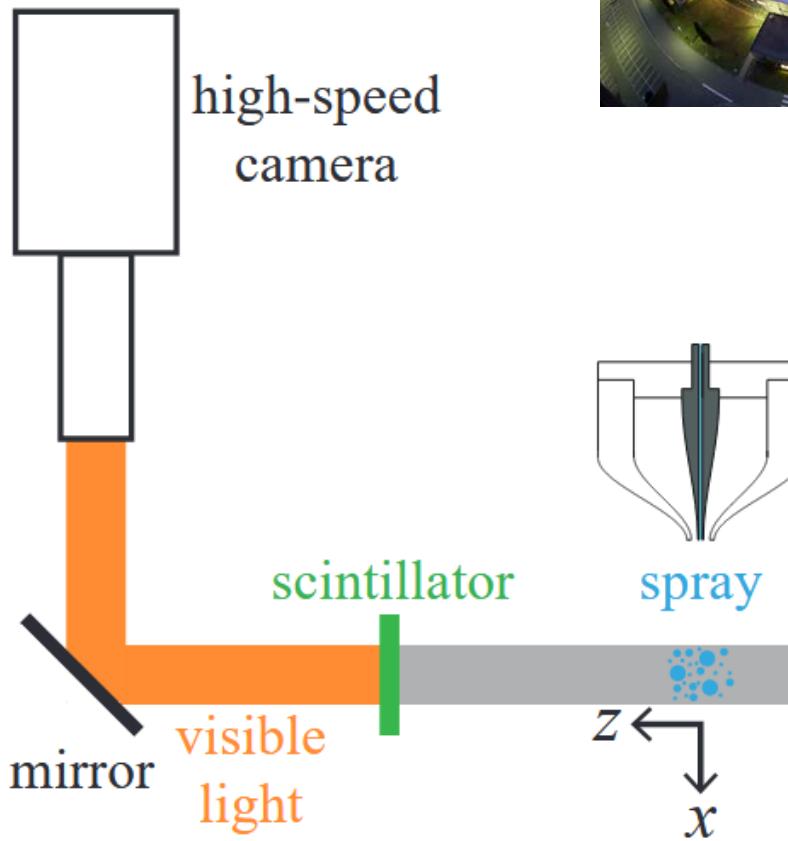


Lasheras & Hopfinger, ARFM 2000

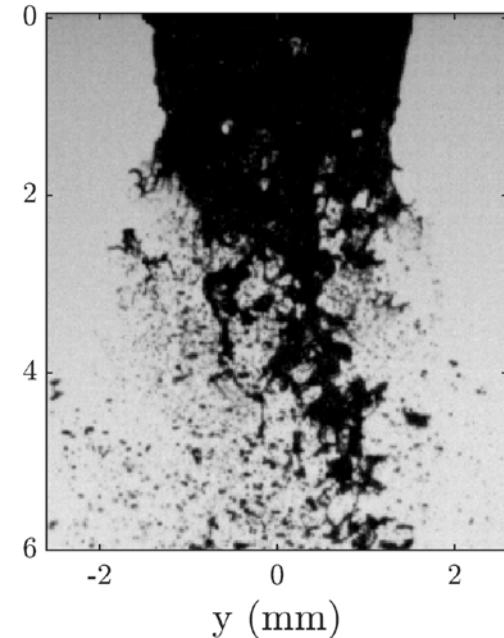
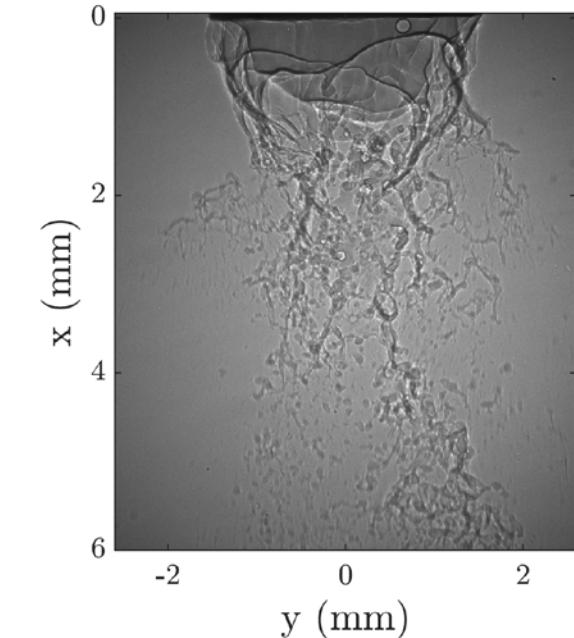




A. Rack
ID19



1 s acquisitions
30 – 100 kHz
7 μm per px
Field of view 7x7 mm 2

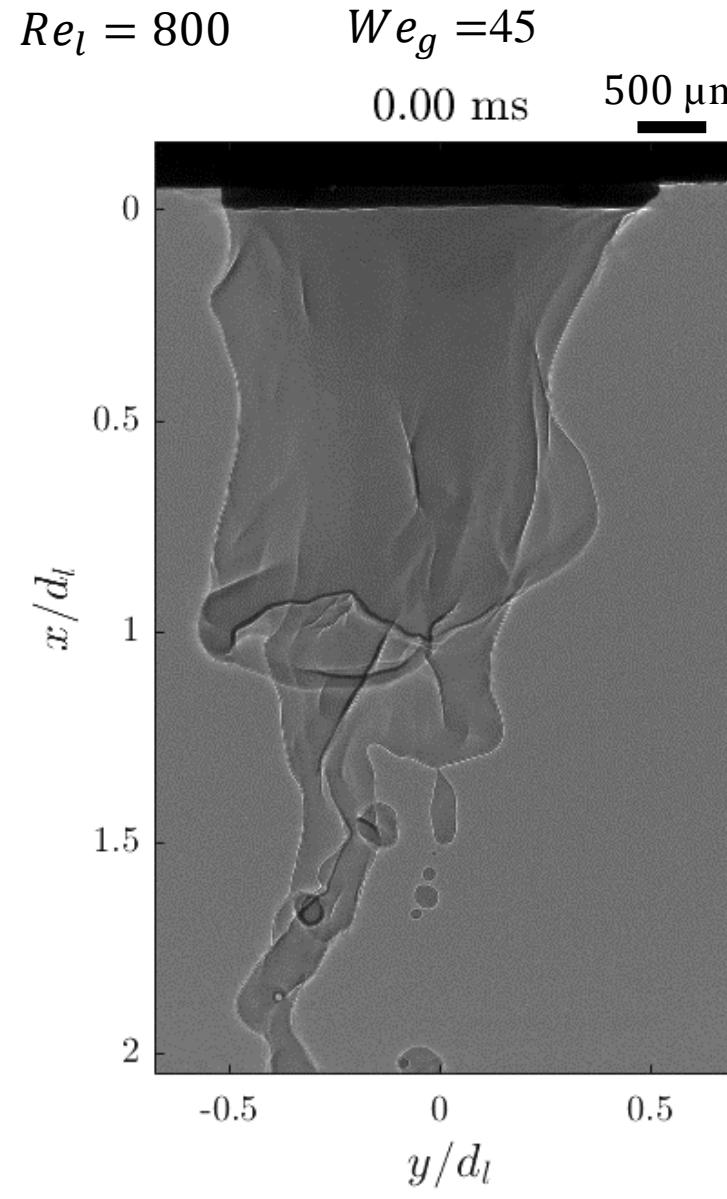


$$800 < Re_l < 5000$$

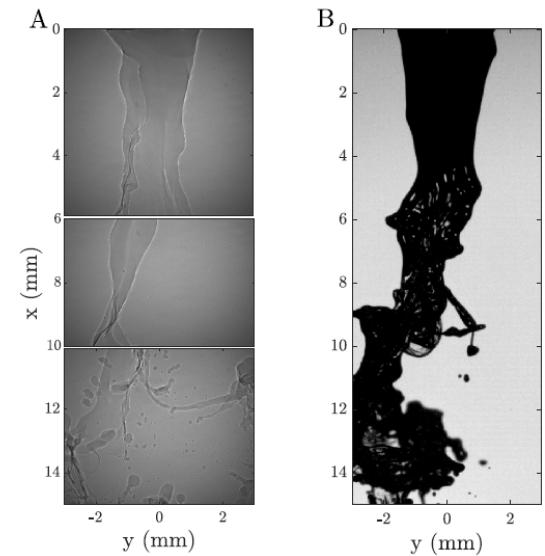
$$2 \cdot 10^4 < Re_g < 2 \cdot 10^5$$

$$30 < We_g < 3200$$

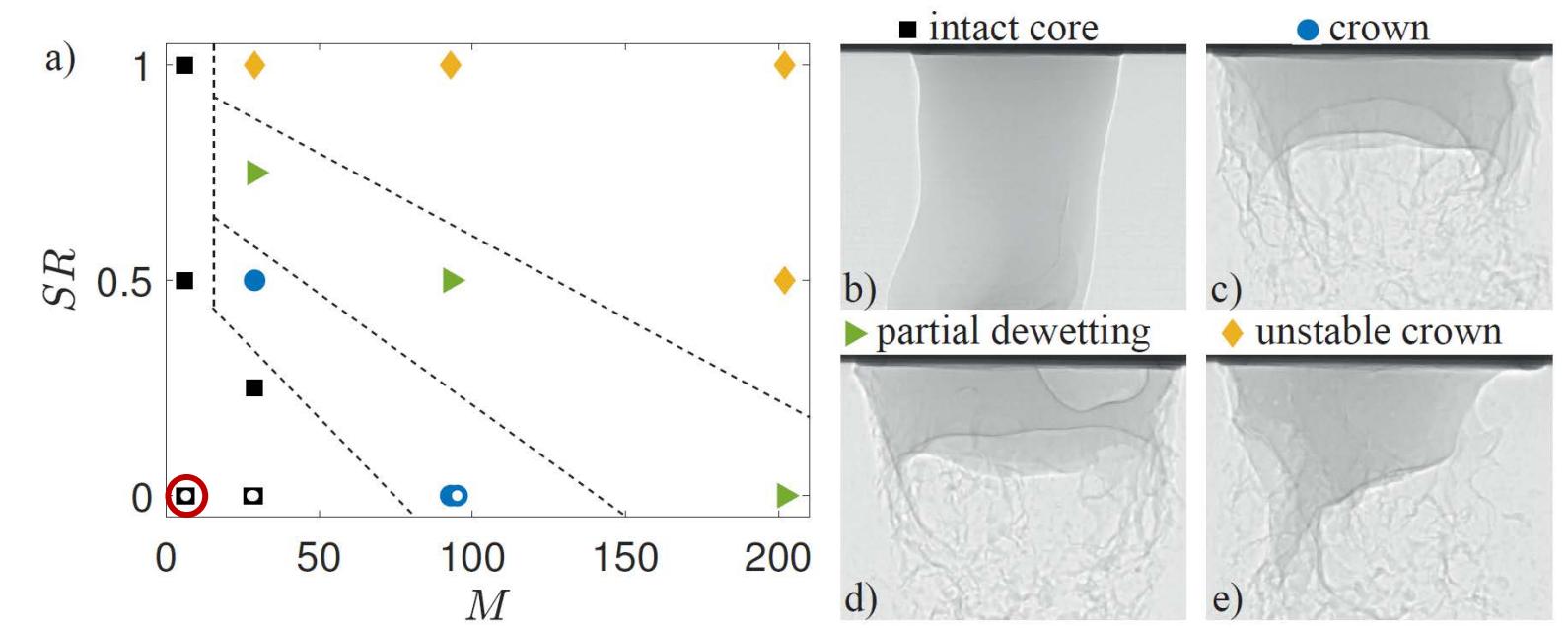
$$5 < M < 700$$



- No gas penetrated within the liquid jet's core
- Strong signature of flapping
- Formation of bags
- Localized Kelvin-Helmholtz perturbations
- Liquid structure reattachments encapsulate large air pockets
- Full wetting of the liquid nozzle with high curvature as the liquid jet is accelerated

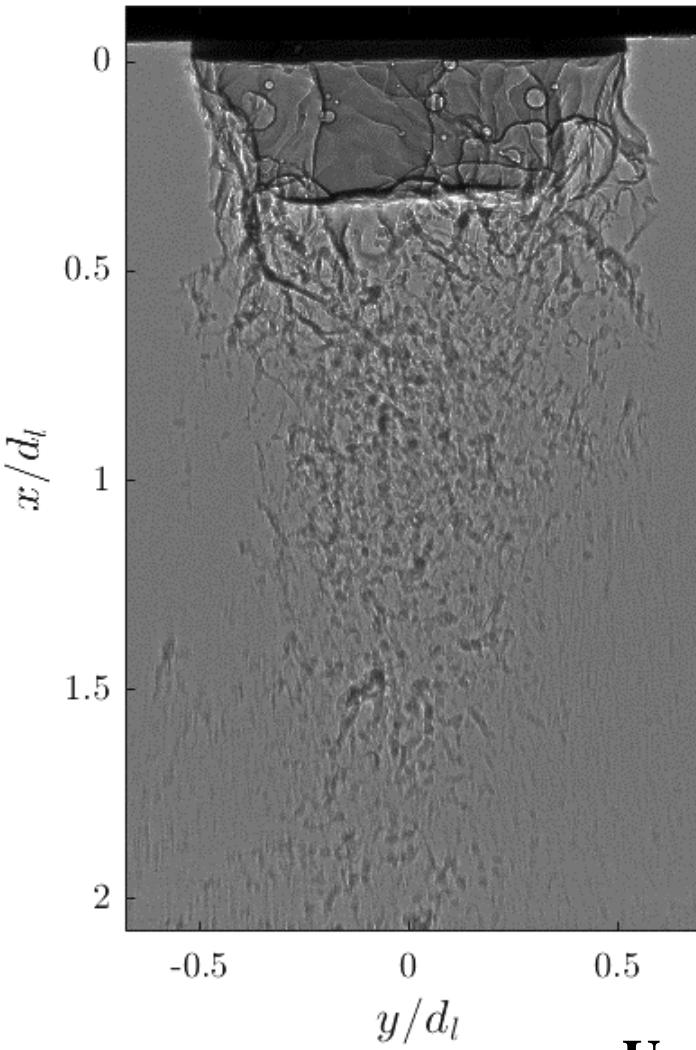


Machicoane et al., IJMF 2019



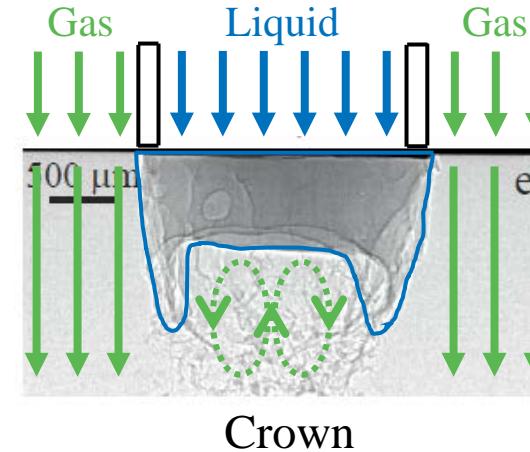
$$Re_l = 1100 \quad We_g = 800$$

0.00 ms 500 μm



Machicoane et al., IJMF 2019

$$Re_l = 1100, We_g = 1350$$



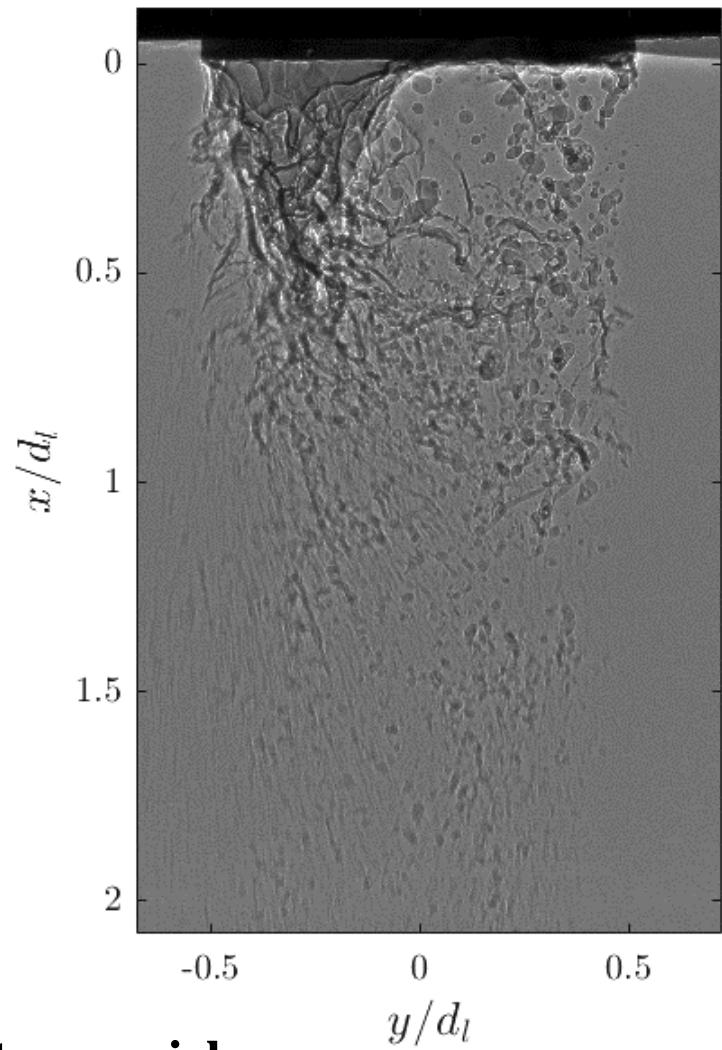
with gas swirl
→ Unstable crown

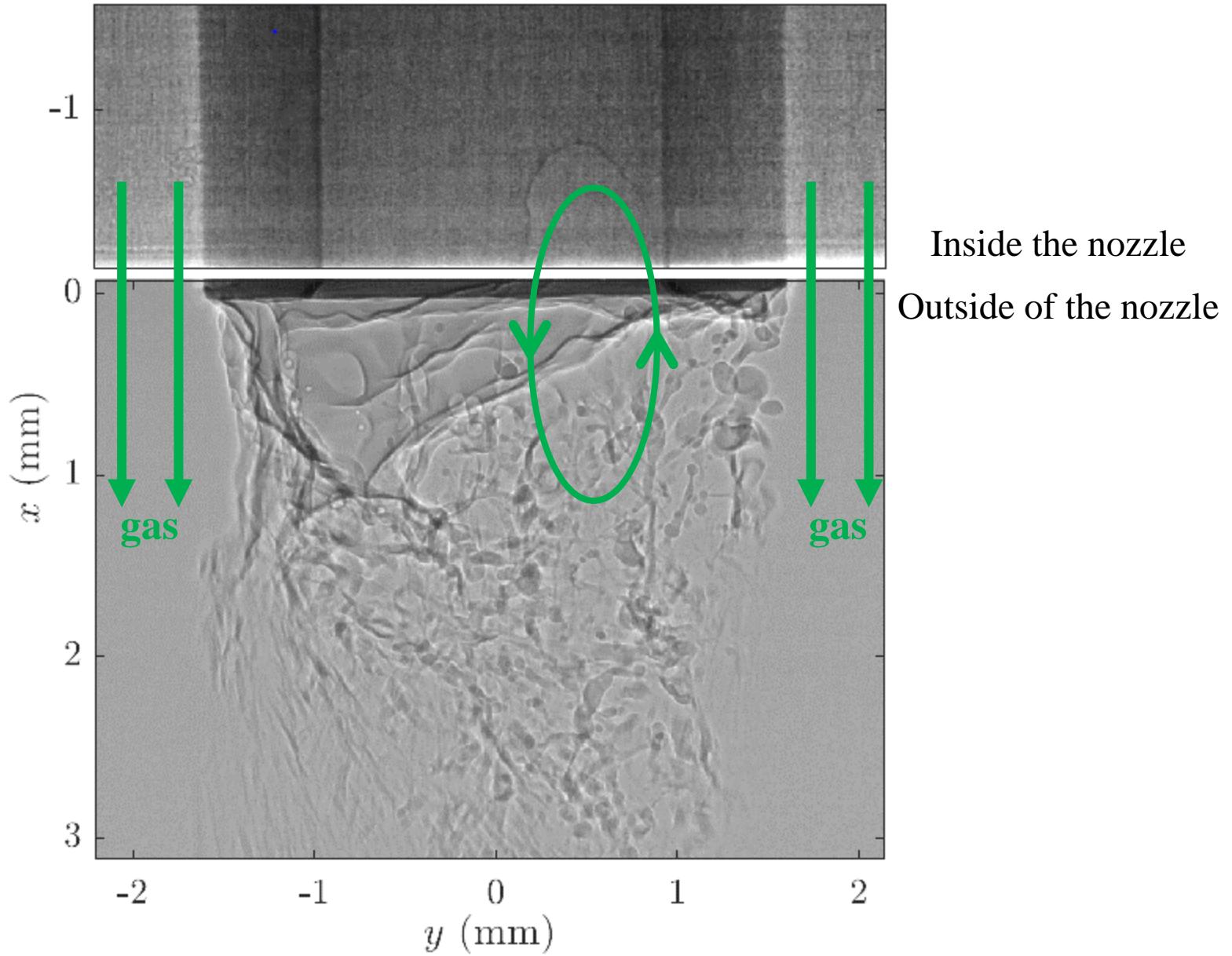


Unstable crown at extreme M values, without gas swirl

$$Re_l = 800 \quad We_g = 1100$$

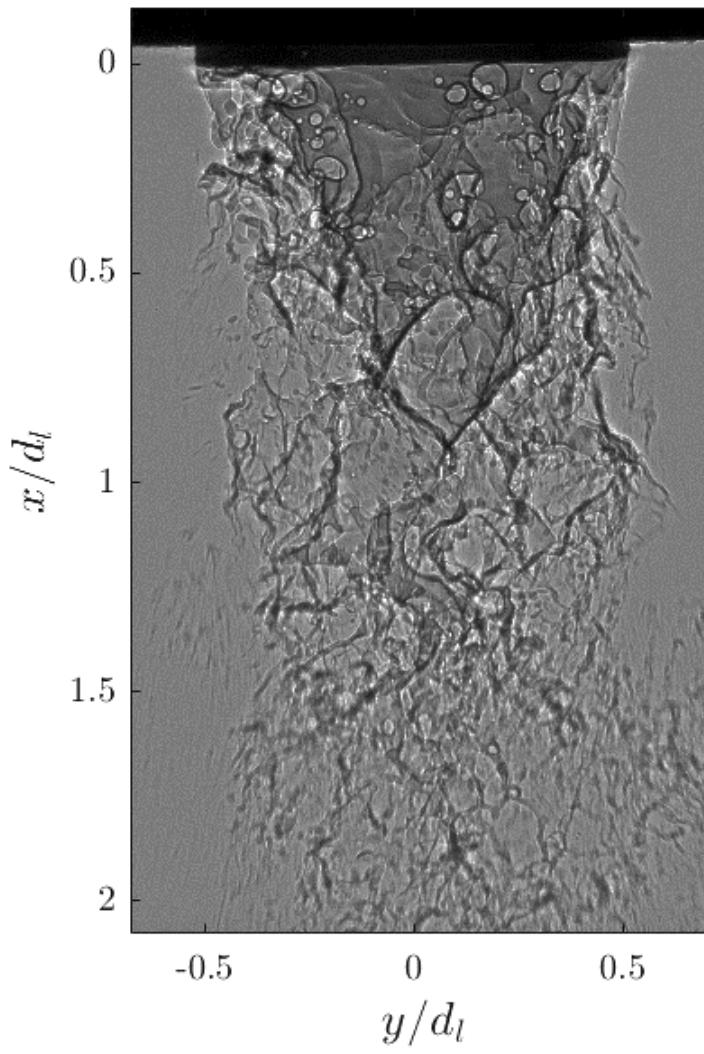
0.00 ms





$$Re_l = 3100 \quad We_g = 800$$

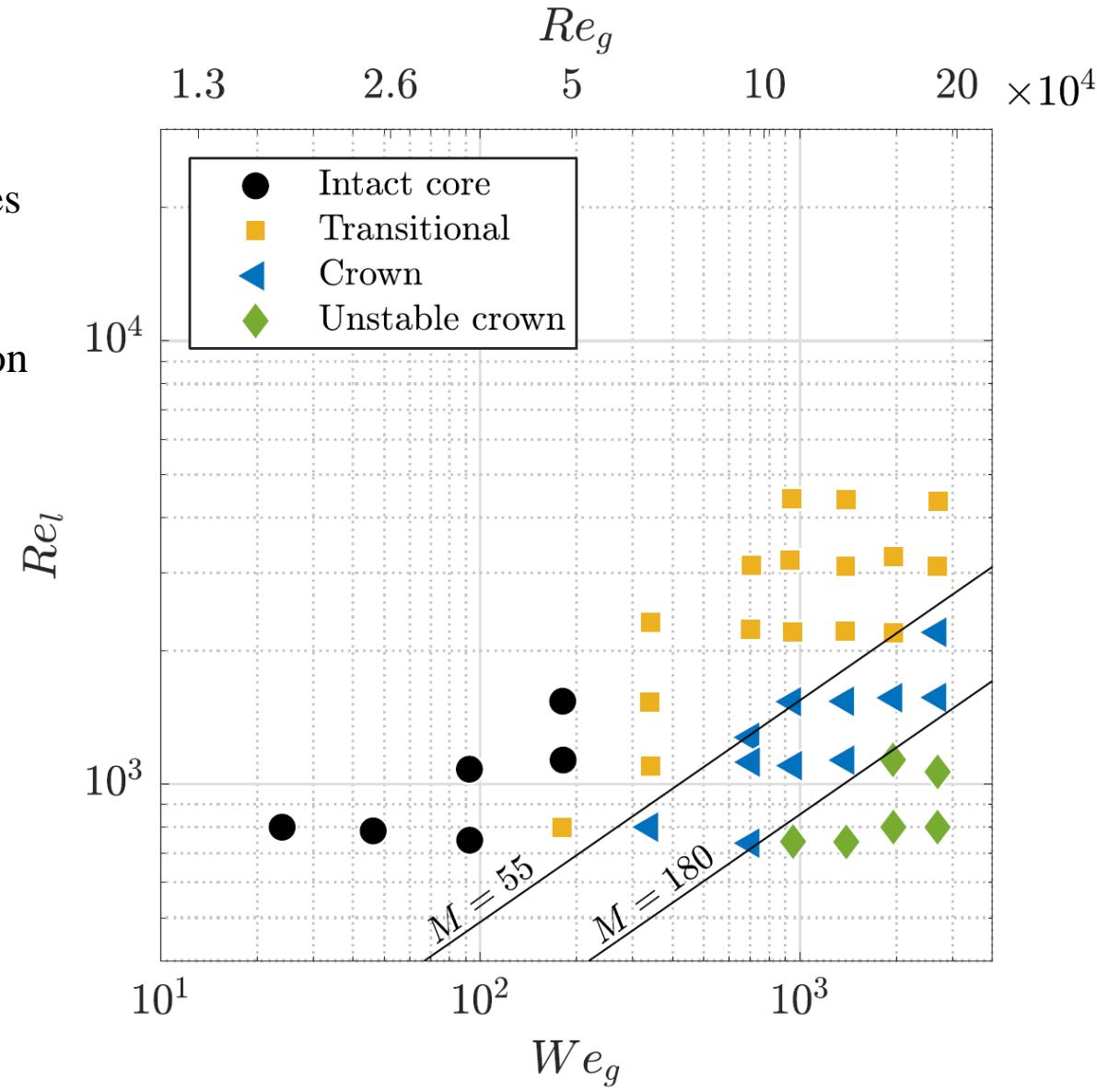
0.00 ms 500 μm

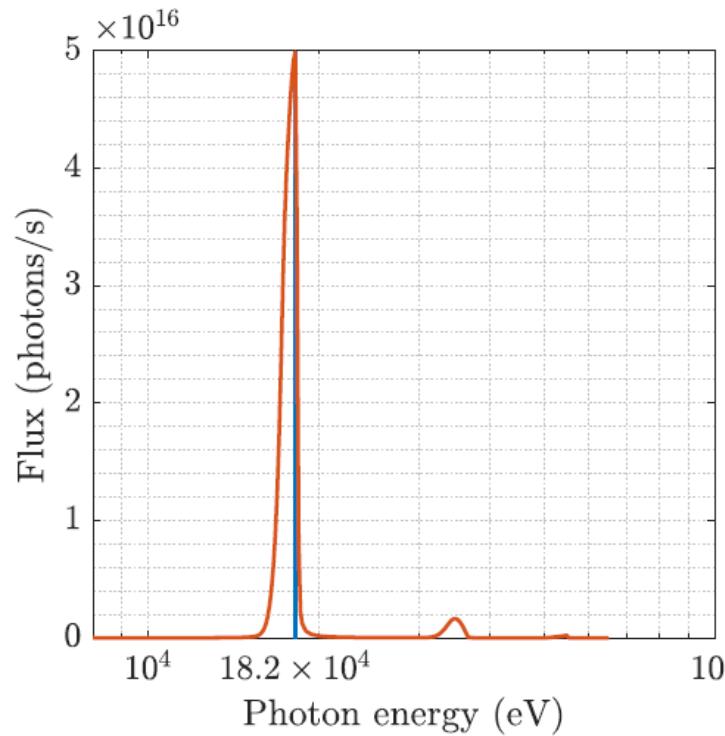


- Many overlapping interfaces
- Many smaller-scale gas recirculations
- « Perforated core » transition between intact core and crown

$$M = \frac{\rho_g u_g^2}{\rho_l u_l^2}$$

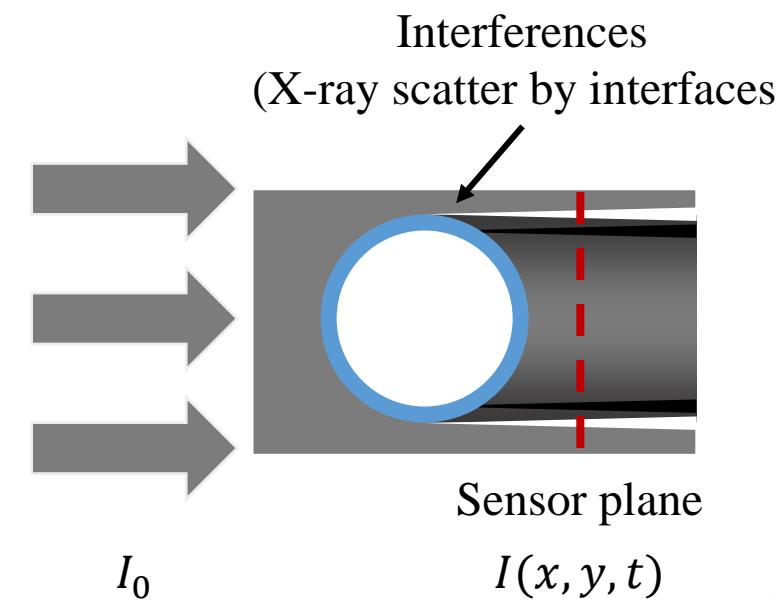
Kinetic energy arguments seem in good agreement for transition to crown and unstable crown





X-ray absorption by the liquid jet follows Beer-Lambert's law

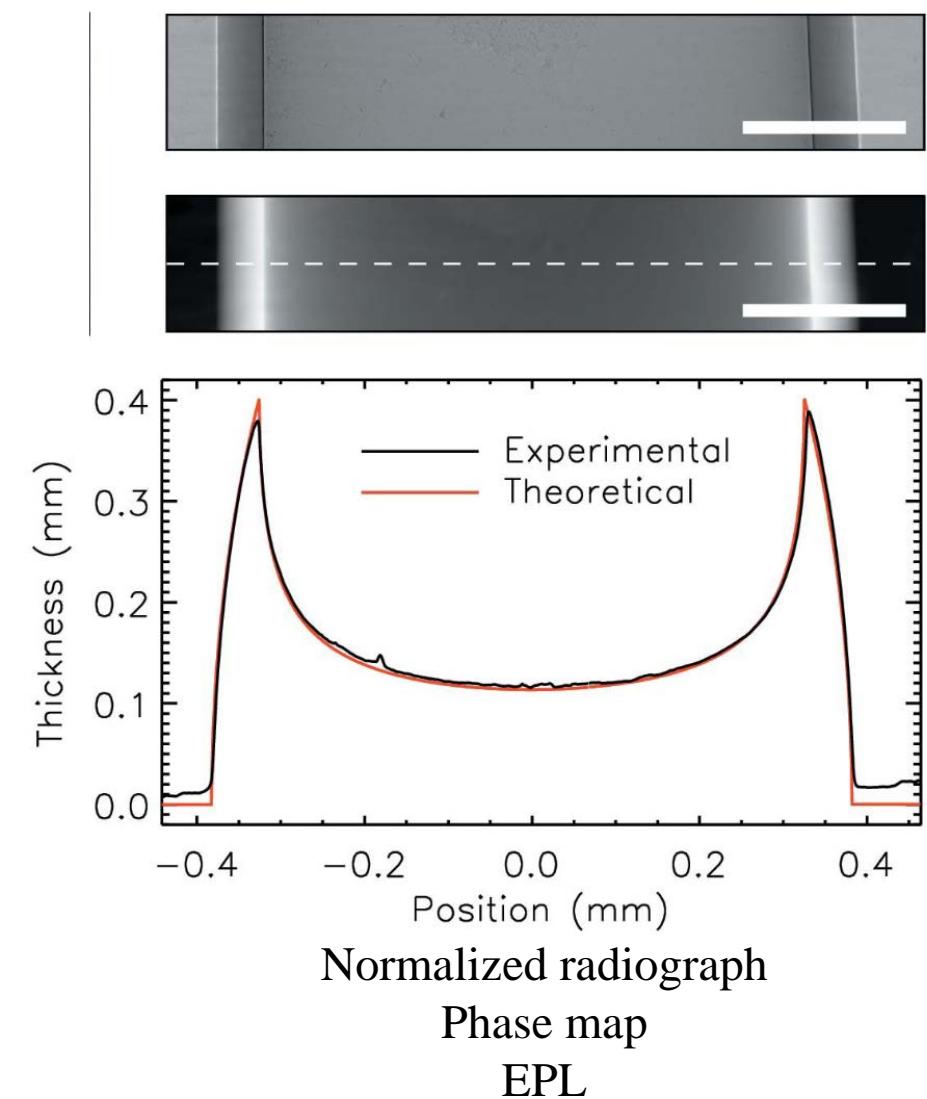
→ Equivalent path length (EPL)

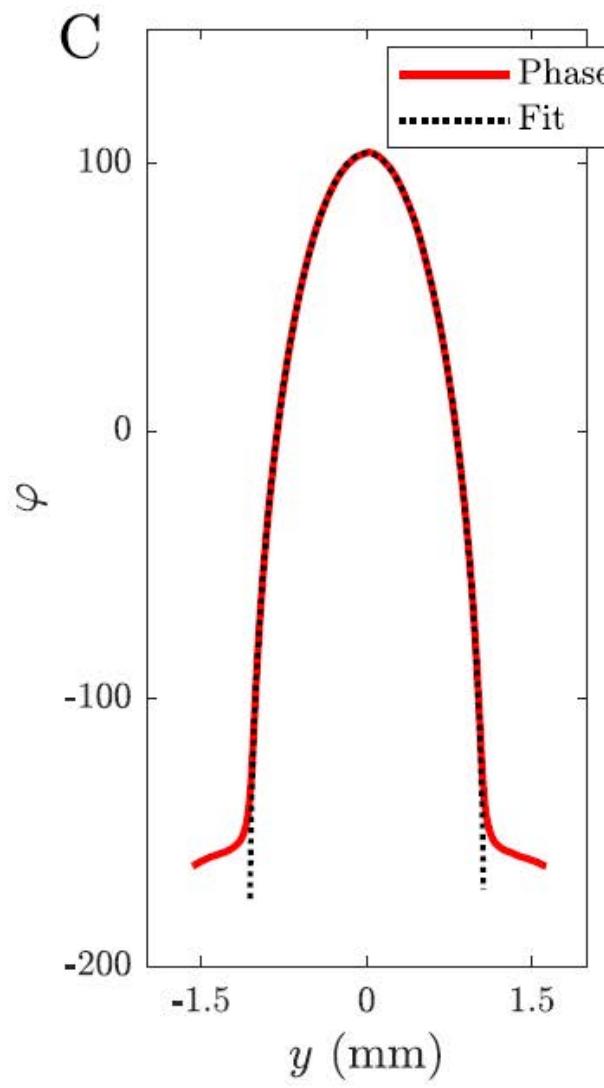
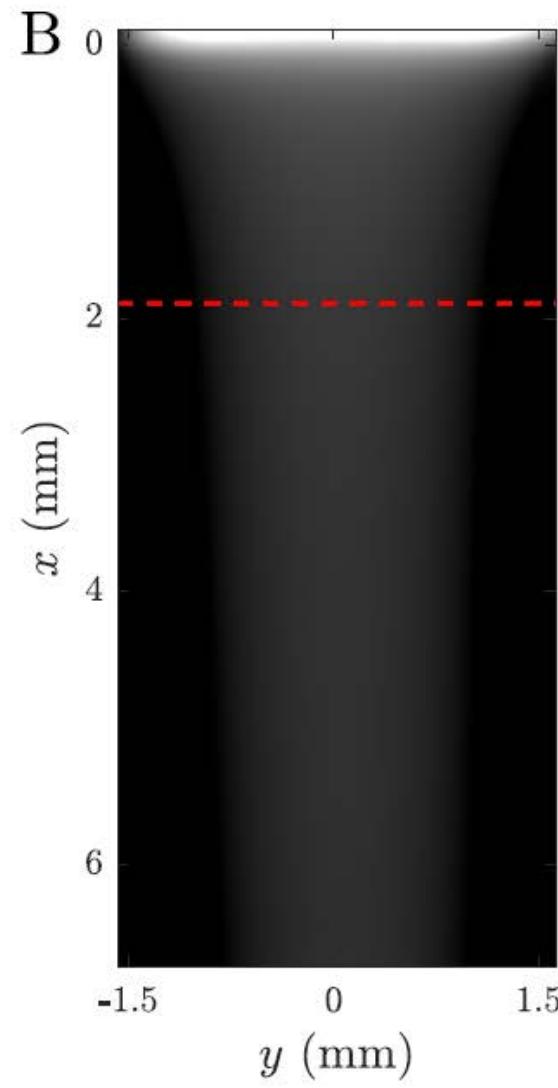
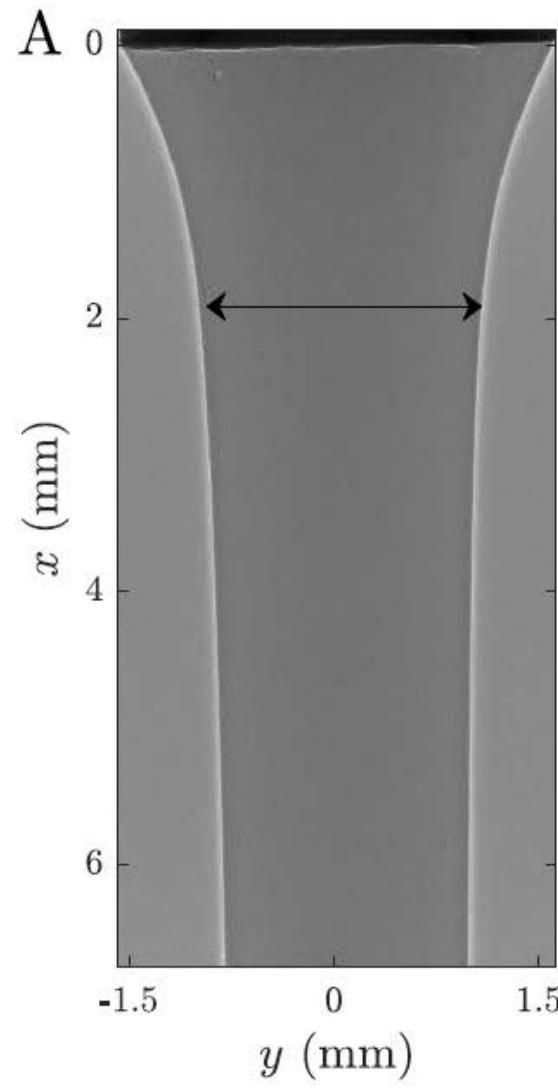


$$\frac{I(x, y, t)}{I_0(x, y)} = \int e^{-\mu(\lambda)h(x, y, t)} d\lambda$$

- Remove interferences $f(x, y, \lambda)$
- Retrieve phase map ϕ
- Convert into EPL map

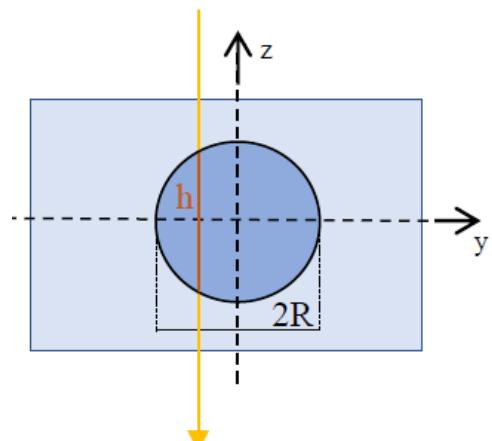
Weitkamp et al.
J. of Synchrotron Radiation 2011





X-ray absorption by the liquid jet follows Beer-Lambert's law

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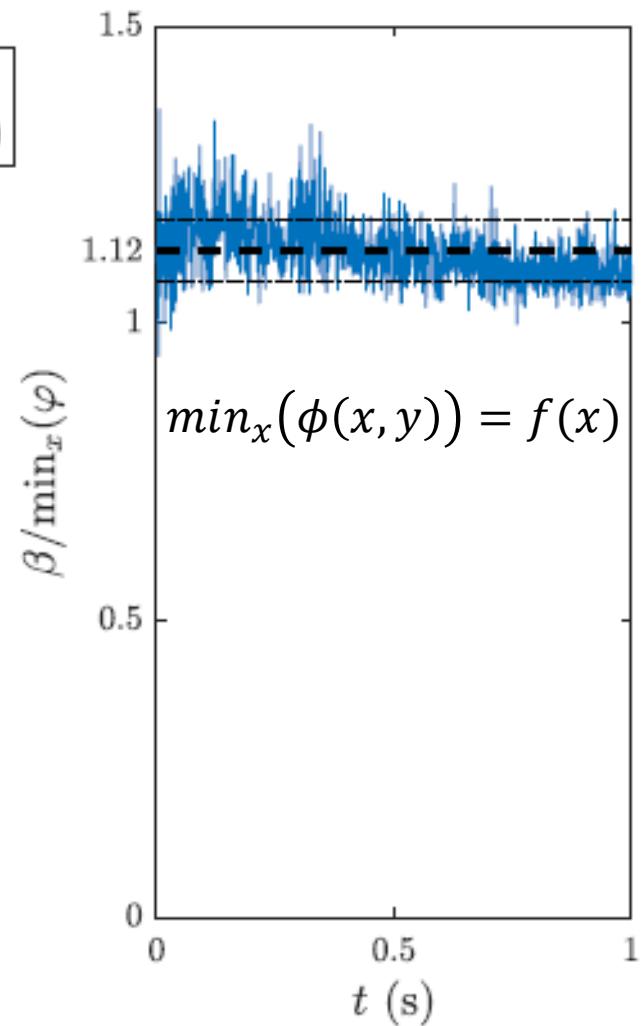
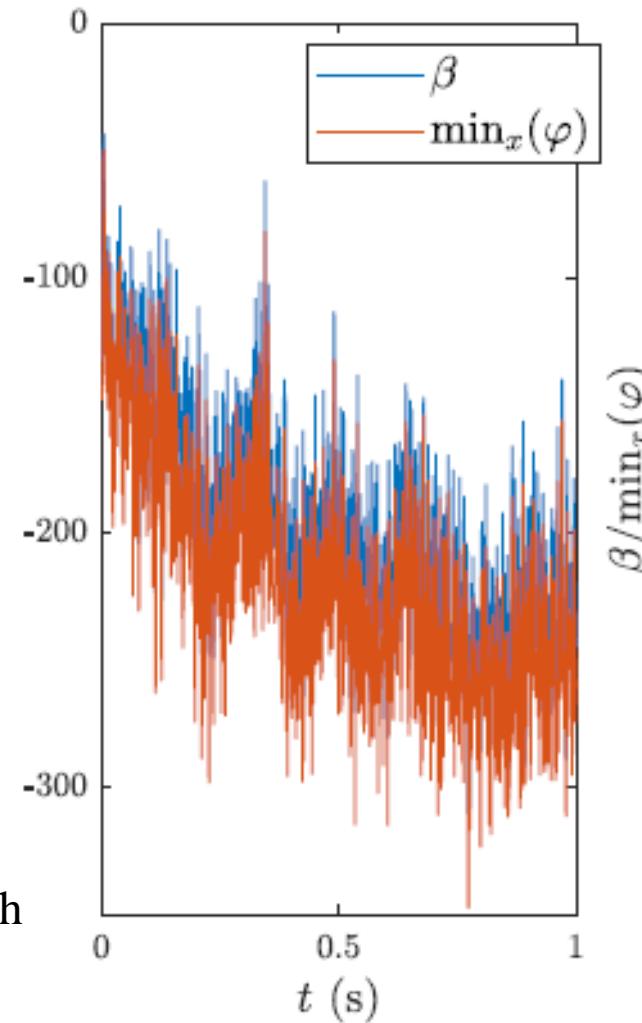
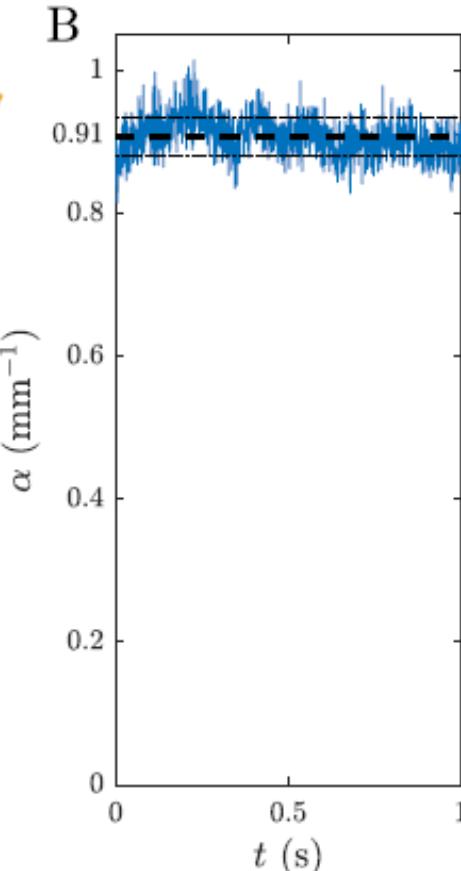
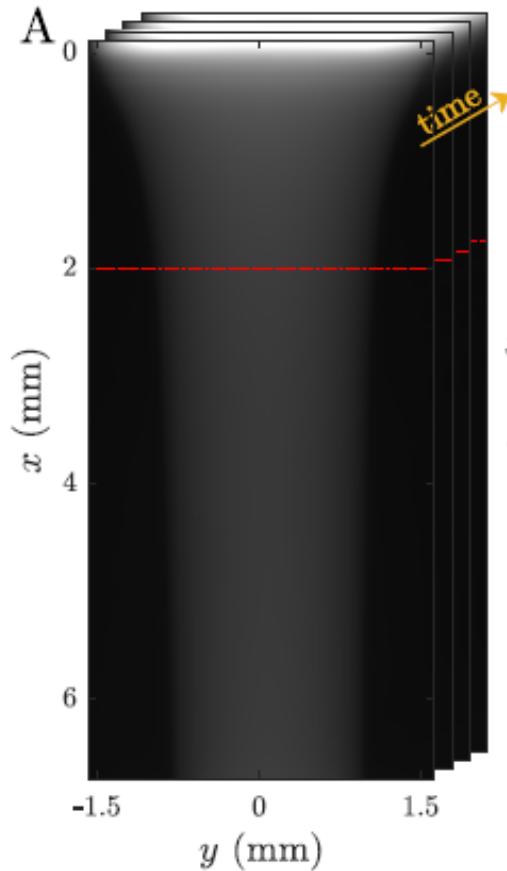


$$\phi = \alpha h(y) + \beta$$

$$h(y) = 2R \sqrt{1 - (y/R)^2}$$

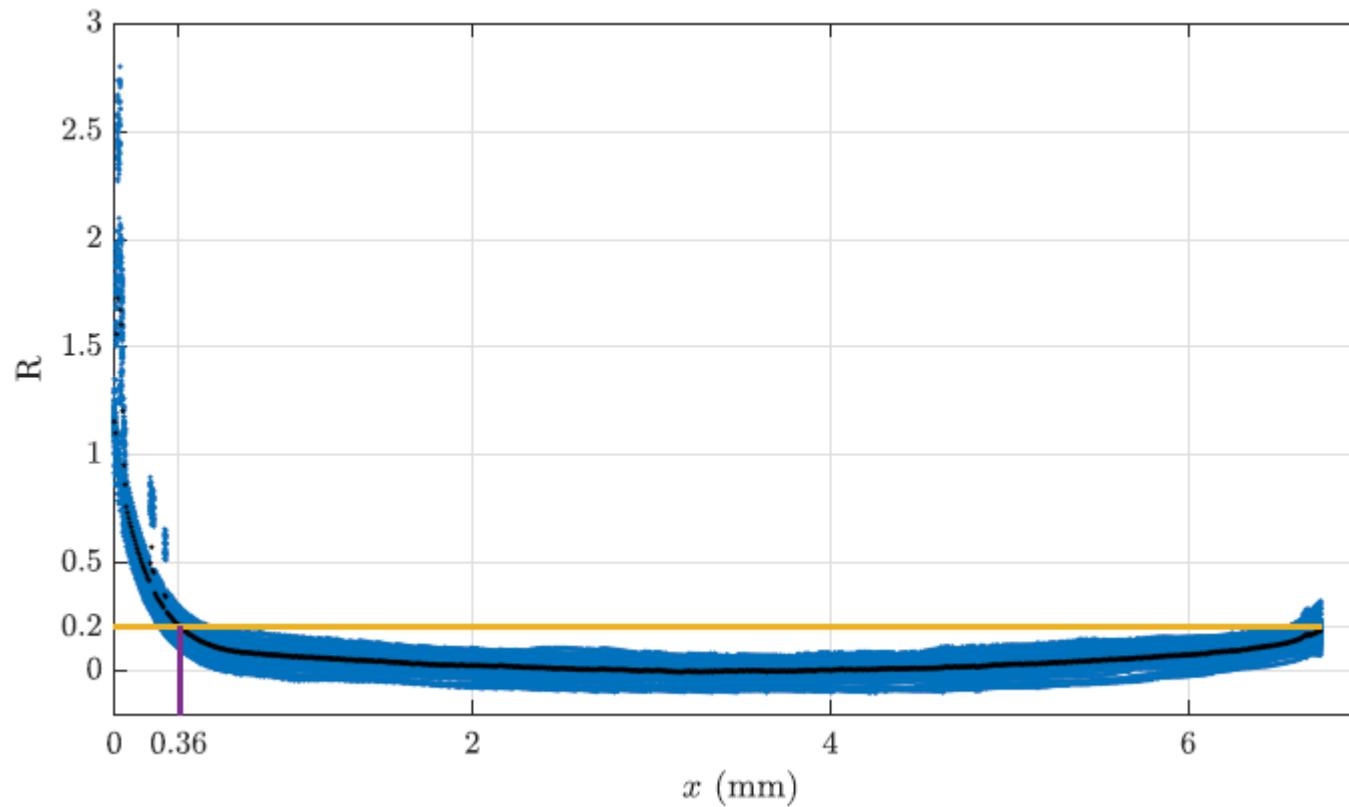
→ Calibrate for the coefficient α and β (non-monochromatic, spatial and temporal inhomogeneities...)

Calibration of the phase maps conversion

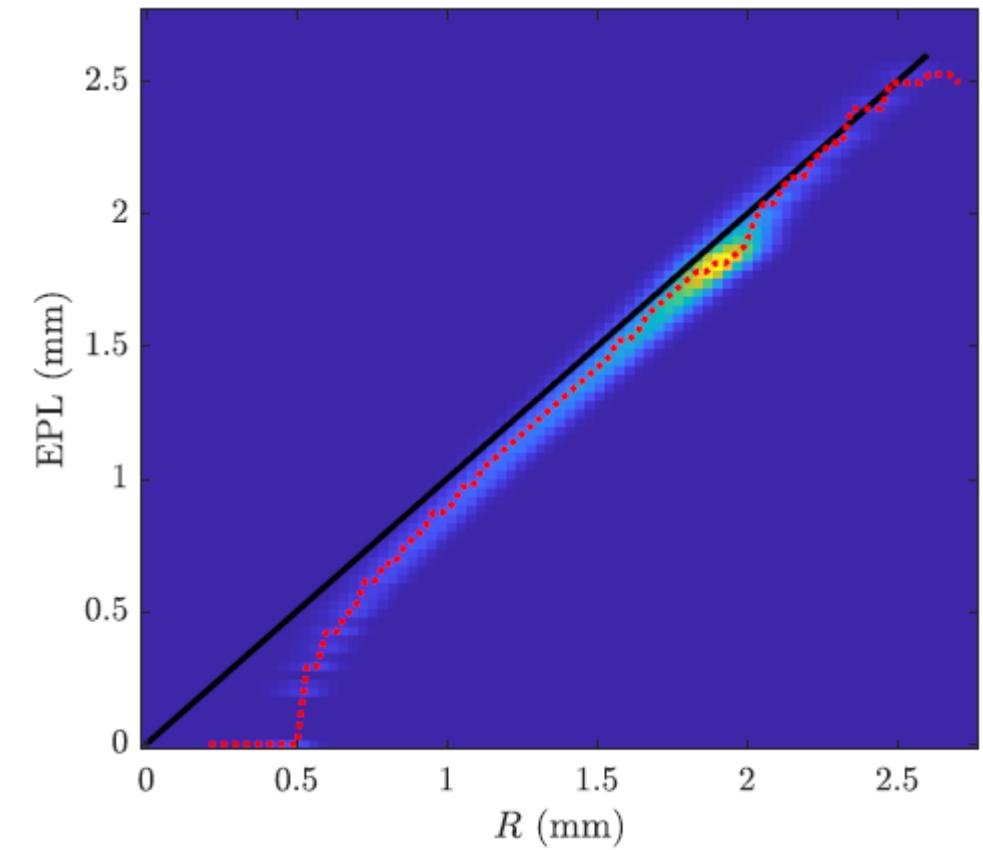


α, β depend on the sample and on the X-ray wavelength

→ $\alpha = 0.91 \text{ mm}^{-1}$ and $\beta = 1.12 \min_x(\phi)$ for what follows



Longitudinal cut along the liquid jet

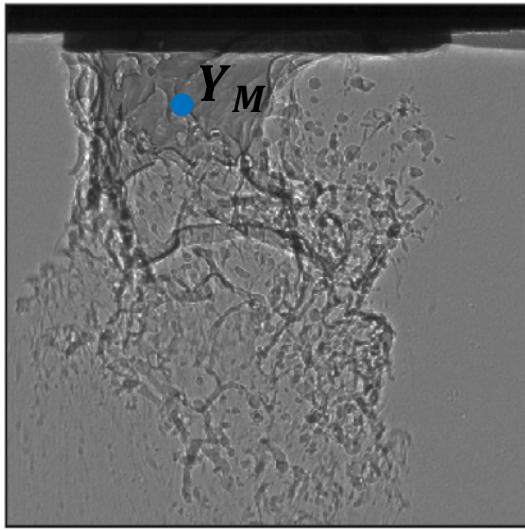


2D PDF: EPL vs fitted radii

Limitations of the uncertainties' evaluation

- Nozzle glare (ANKA Phase is for a single material)
- Interference patterns due to X-ray scattering by interfaces limit the probing of small radius values
- ➔ For $x > \frac{D_l}{10}$ and for $EPL > 1$ mm, approximately 10% accuracy ($\sim 20\%$ for smaller thicknesses?)

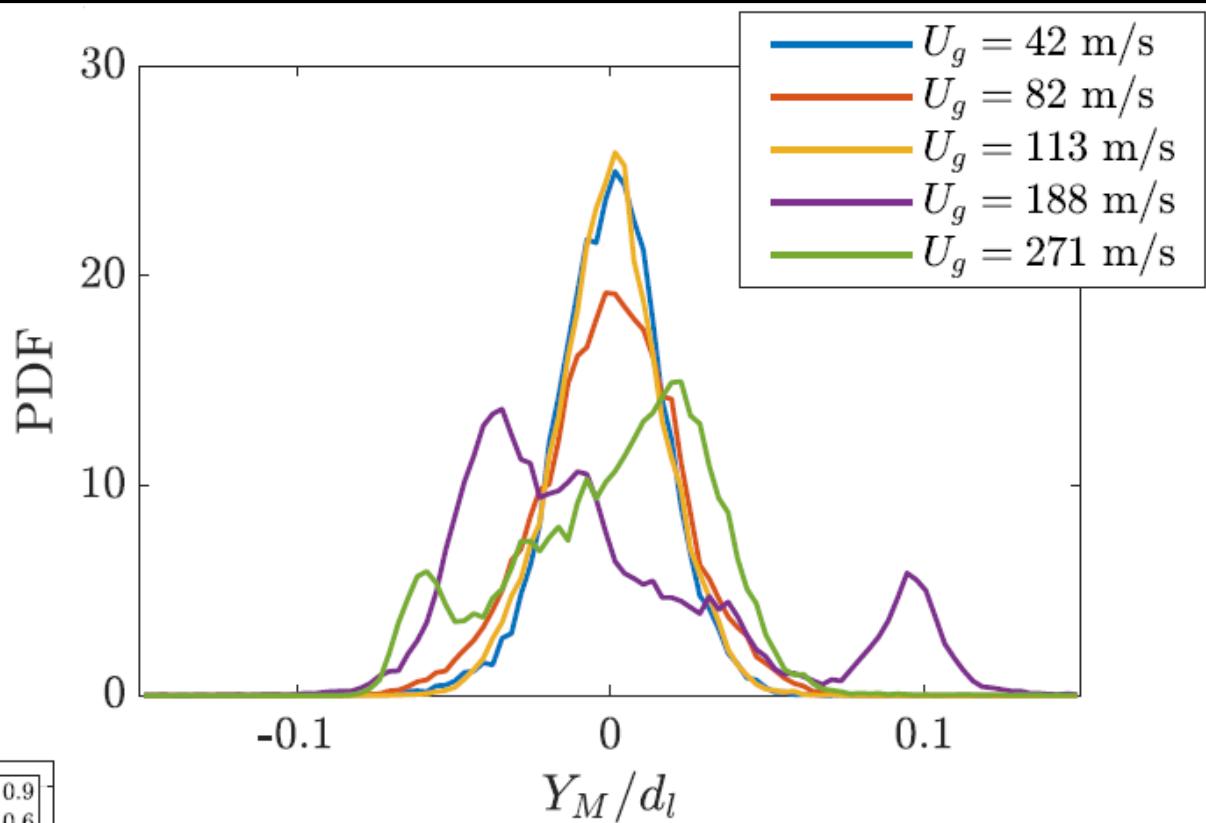
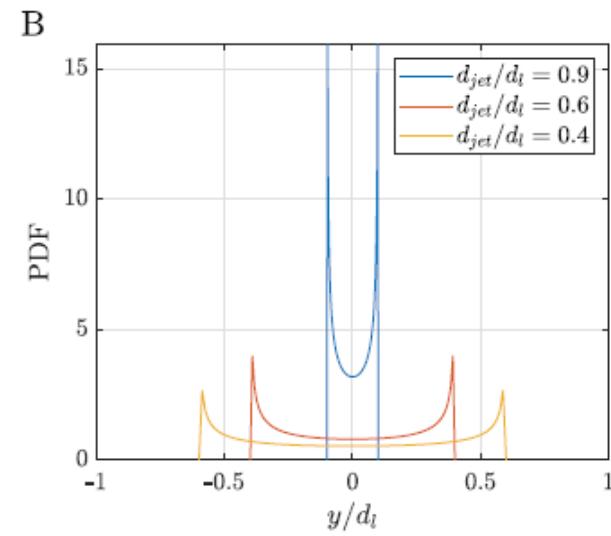
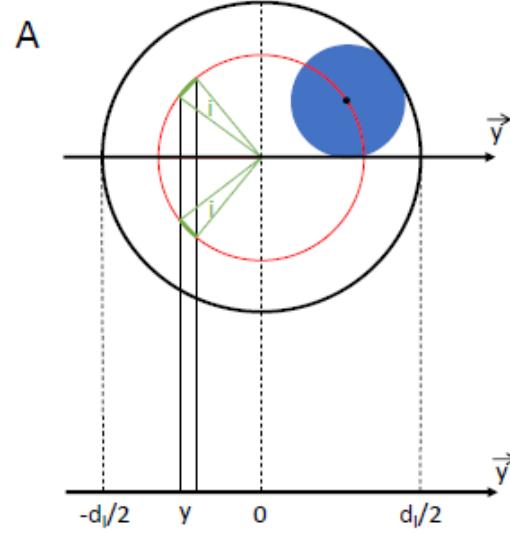
Transverse center of mass of the liquid core



$Re_l = 800$

$40 < We_g < 2000$

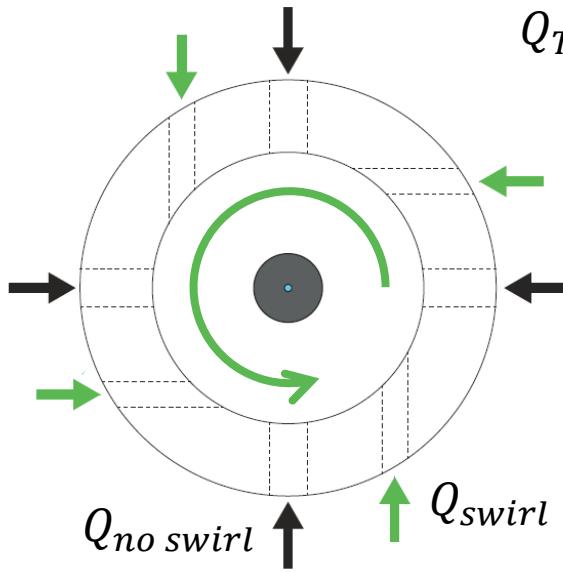
Center of mass along y



Simple model of the unstable liquid crown

- Circular cross-section
- Anchored at a given distance from the nozzle's axis d_{jet} that is varied
- Qualitatively explains the shape of the PDF

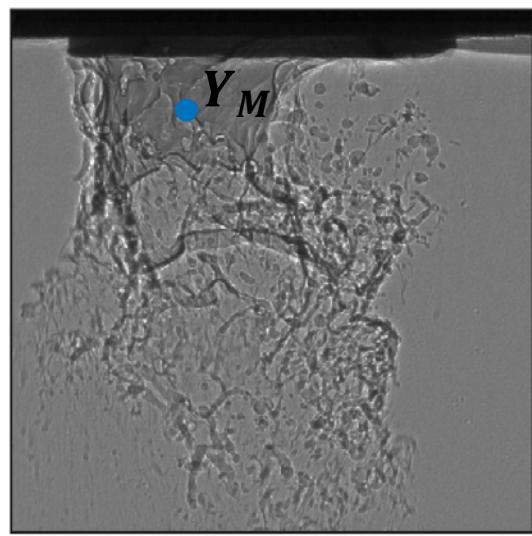
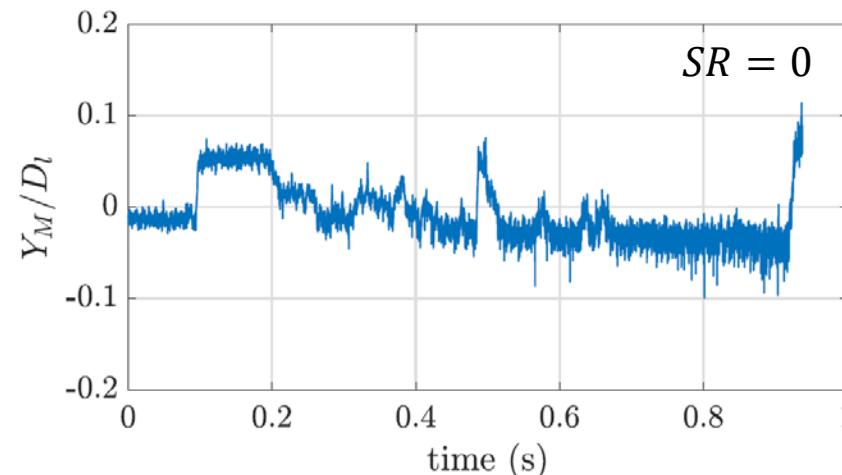
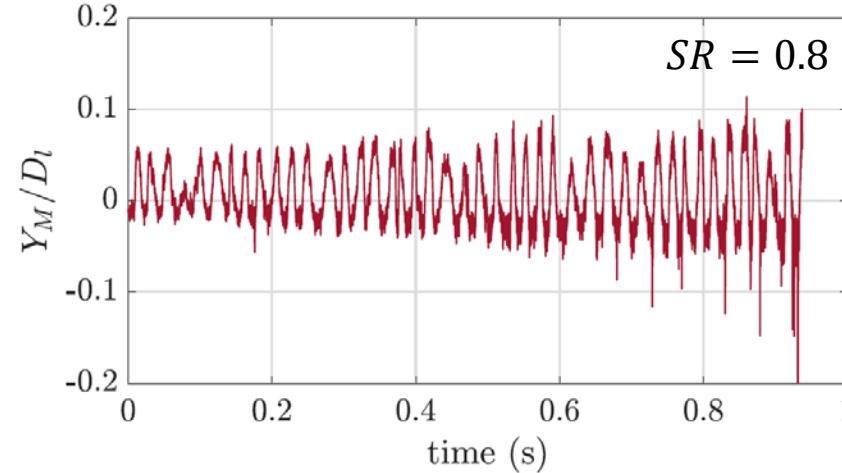
Temporal dynamics of the unstable crown and role of swirl



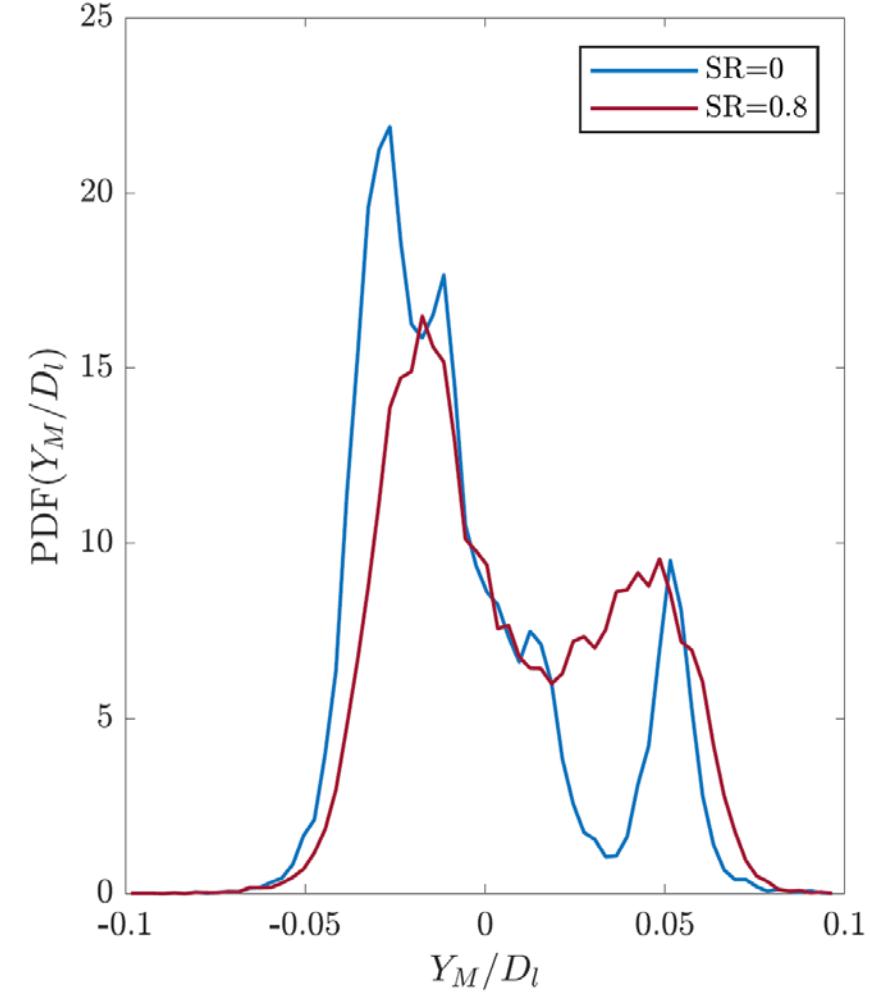
$$Q_{Total} = Q_{SW} + Q_{NS} = cst$$

$$SR = Q_{swirl}/Q_{no\ swirl}$$

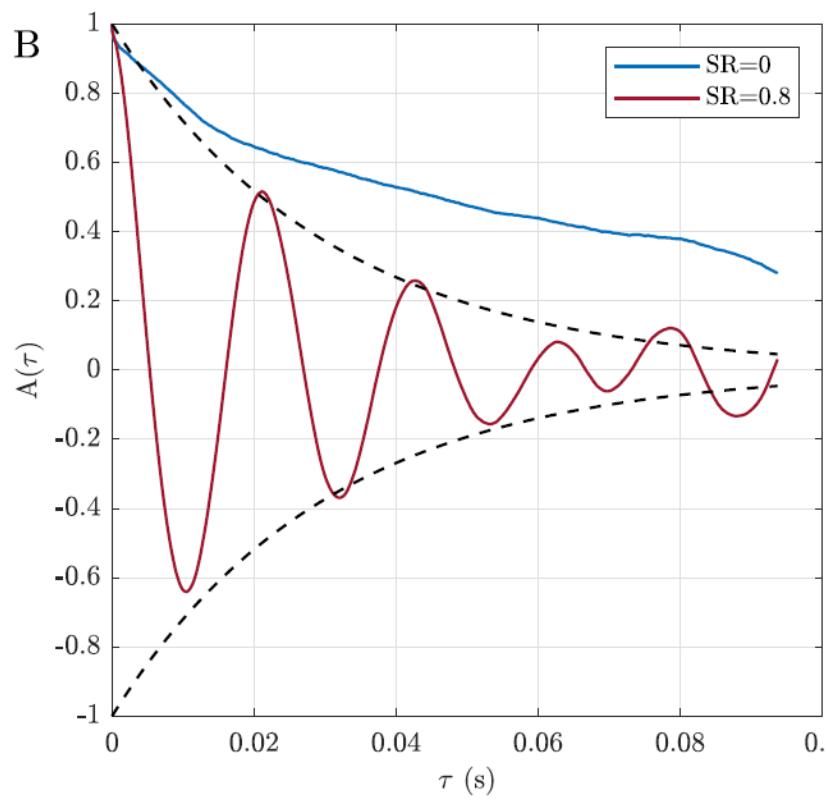
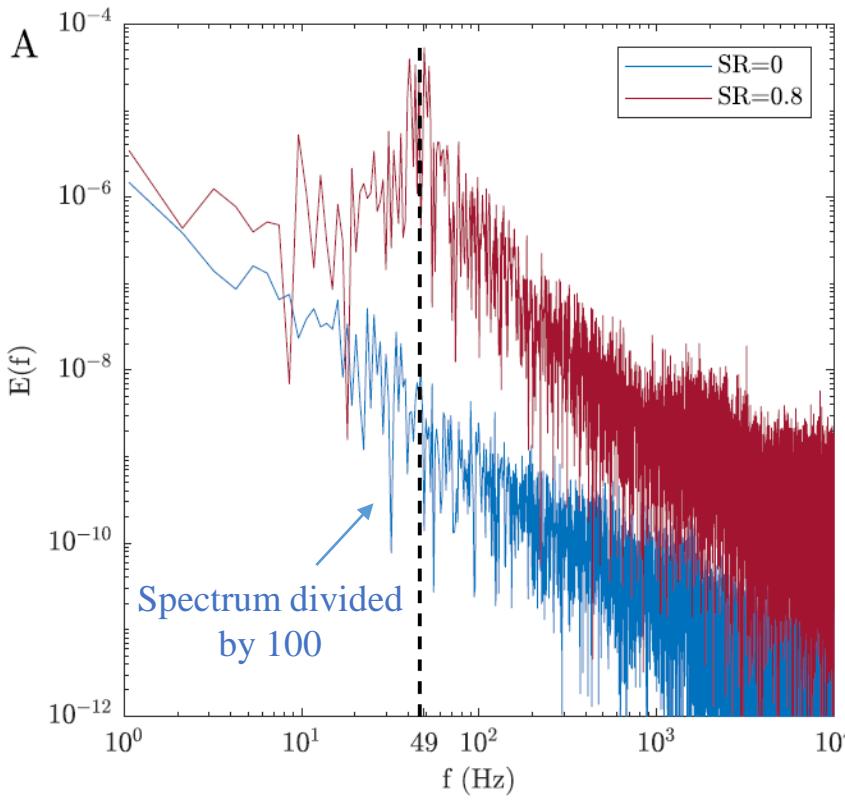
$$Re_l = 800 \quad We_g = 950$$



Center of mass along y



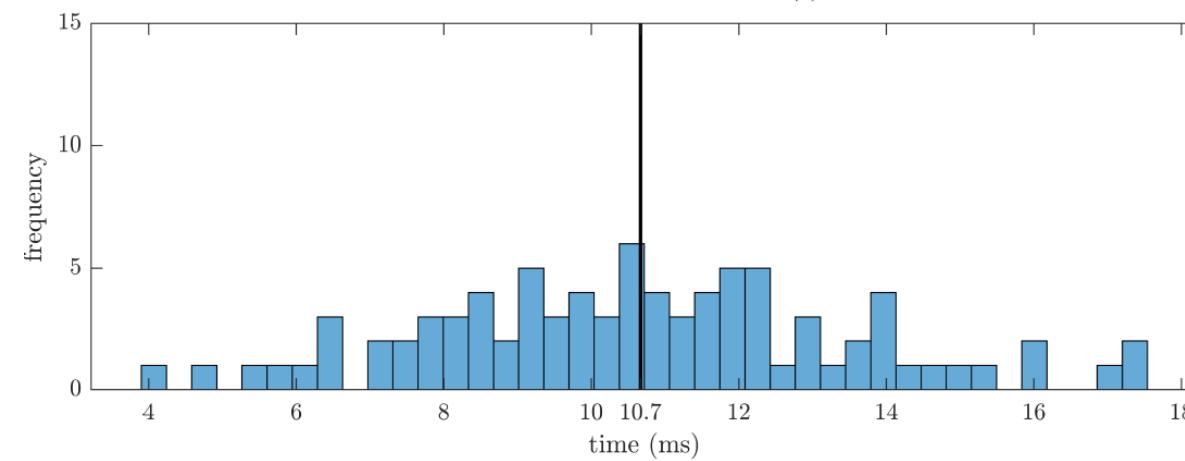
Temporal dynamic of the unstable crown



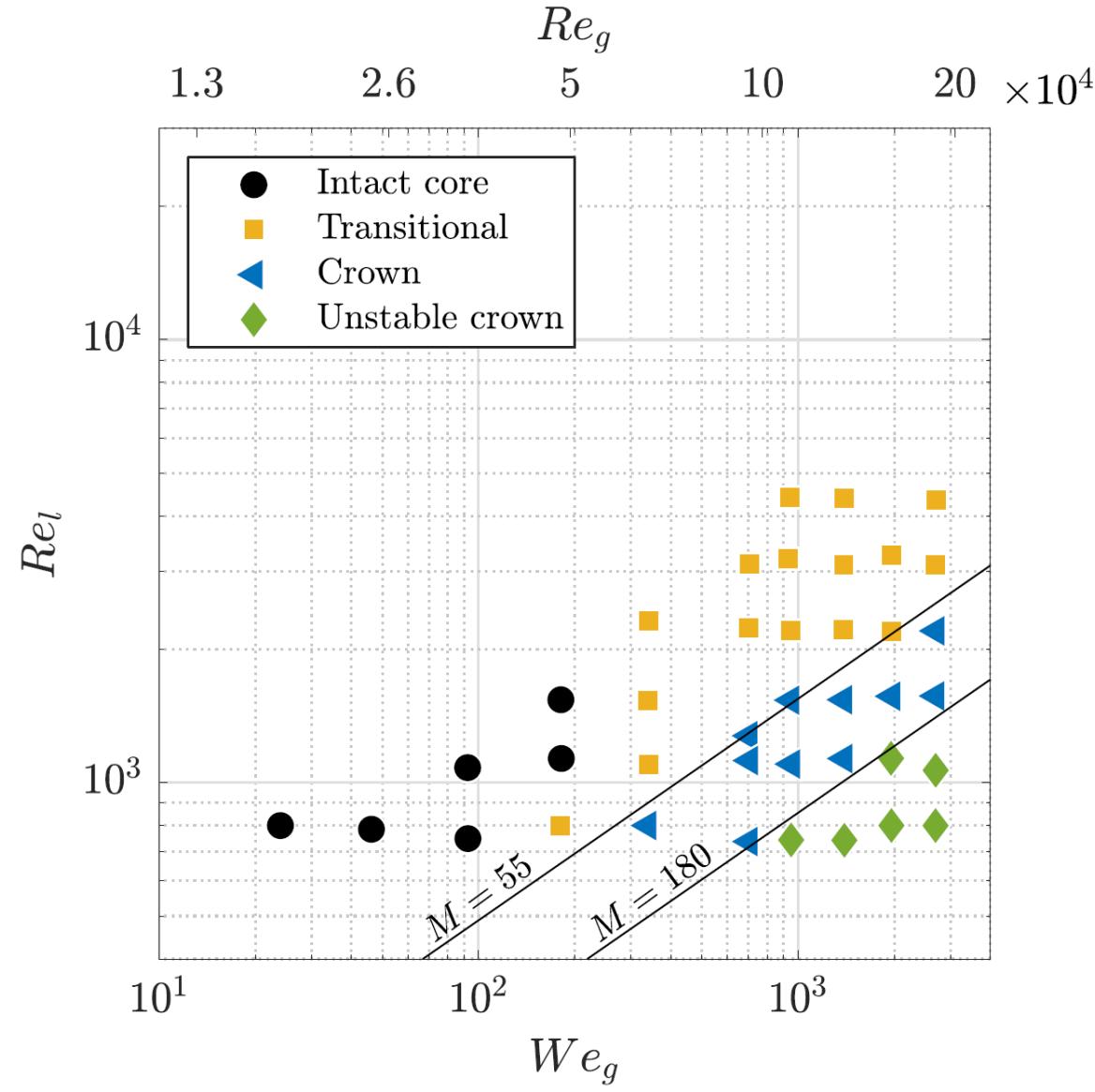
$$Re_l = 800 \quad We_g = 950$$

- Strong periodicity signature with swirl
- Mean residence time on the side is twice the oscillation period
- Without swirl, onset of a slow dynamic which would required longer acquisition to investigate
- Decorrelation of the liquid core center of mass is orders of magnitude slower than that of the liquid core length

SR = 0.8
 Bin the signal into « left », « right », and « center »
 → Histogram of the residence times spent « on the side »



- Proposed a method to retrieve liquid path using X-ray
- At higher We_g , liquid core undergoes transitions, up to unstable crown, even without gas swirl
 - Intact liquid core
 - Transitional liquid core
 - Liquid crown
 - Unstable liquid crown
- Gas swirl leads to
 - Earlier onset of unstable crown (i.e. at lower We_g)
 - Much more frequent motions of the gas recirculation
 - Similar PDF for center of mass
- Open questions
 - Regime map (Re_l , We_g , SR)
 - Characteristic frequency of the liquid core motions with and without swirl



(see *Kaczmarek et al., IJMF 2022* for flapping and role of swirl at lower We_g)